

Real-time estimation of kinematic variables for advanced mechanical systems

[Workshop] Dynamics and control of classical and quantum
systems

Nicolas Petit

Centre Automatique et Systèmes (CAS)
Mines Paris, Université PSL

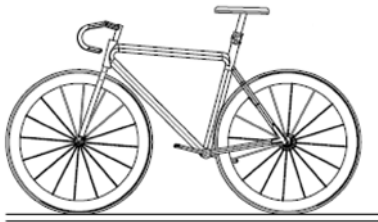
June 2, 2026



Mechanical systems

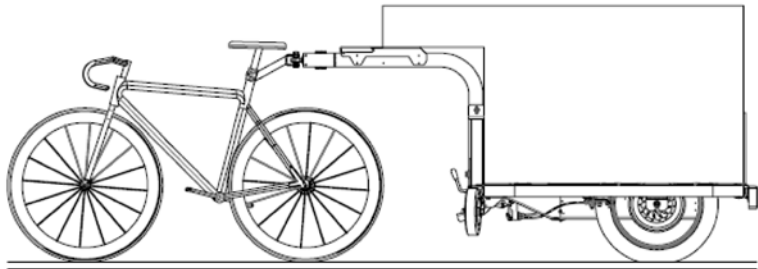
- Essential quantities of mechanical systems include **generalized coordinates** (positions, velocities, angles, angular velocities, *motional variables* for non-holonomic constraints), inertia, as well as internal and external **forces**
- **Mastering the motion** to carefully selected **reference signals**
- Typical instruments include **accelerometers, gyroscopes, strain gauges**, radar, magnetometers, resonators, Pitot tubes, PIV (Particle Image Velocimetry), ...
- **Core claim**: stabilize the unstable, accelerate the sluggish, unburden the heavy

Dynamics transformation by feedback on internal forces



Operates on a **simple open-loop principle**: a sensor in crankset detects rotation, triggers the hub motor, monitors rotation speed and contributes to propulsion based on look-up tables (“Eco” / “Comfort” / “Turbo”)

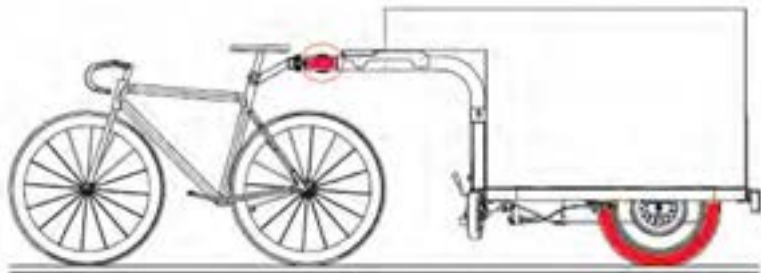
Dynamics transformation by feedback on internal forces



Attach a heavy trailer (100-500 Kg) and try to make it **transparent**

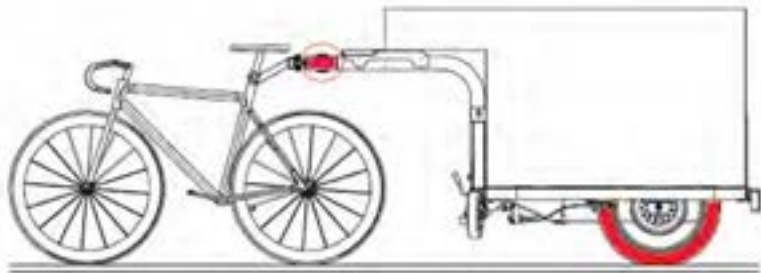
The **basic strategy is not sufficient**. “Turbo” on all the time.

Dynamics transformation by feedback on internal forces



Trailer has powerful **motors**, mechanical interaction occurs **only** at the **drawbar**

Dynamics transformation by feedback on internal forces

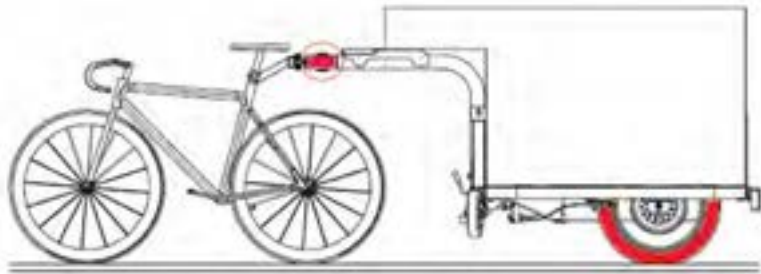


$$\dot{x} = f(x, \underbrace{u_1}_{\text{cyclist}}, \underbrace{u_2}_{\text{electric drives}}, \underbrace{p}_{\text{(masses, geometry, ...)}}), \quad y = h(x, u, p)$$

Model assignment - “dynamics shaping”

$$\begin{cases} \dot{x} = f(x, u_1, k(\xi), p) \approx f(x, u_1, 0, p_0) : \text{bike only} \\ \dot{\xi} = g(\xi, y) \end{cases}$$

Dynamics transformation by feedback on internal forces



$$\dot{x} = f(x, \underbrace{u_1}_{\text{cyclist}}, \underbrace{u_2}_{\text{electric drives}}, \underbrace{p}_{\text{(masses, geometry, ...)}}), \quad y = h(x, u, p)$$

Model assignment - “dynamics shaping”

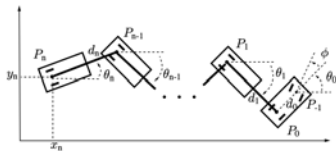
$$\begin{cases} \dot{x} = f(x, u_1, k(\xi), p) \approx f(x, u_1, 0, p_0) : \text{bike only} \\ \dot{\xi} = g(\xi, y) \end{cases}$$

Mass cancellation is achieved

K-Ryole in Paris (10 years of R&D, 1500 vehic.)



The n trailers



Flatness and motion planning : the car with n trailers. *

P. ROUCHON¹ M. FLEISS² J. LÉVINE³ P. MARTIN⁴

Abstract A solution of the motion planning without obstacles for the nonholonomic system describing a car with n trailers is proposed. This solution relies basically on the fact that the system is flat with the cartesian coordinates of the last trailer as linearizing output. The Frobenius formulae are used to simplify the calculation. The 2-trailers case is treated in details and illustrated through parking simulations.

Key words Nonholonomic motion planning, dynamic feedback linearization, linearizing output, flatness, mobile robots.

1992 paper, a lasting inspiration

Inertial measurement units (IMUs), multiple axis



Safran HRG



SBG



Bosch BNO055

Weaknesses: vibrational/acoustic disturbances, converter flips ...

Prices: $\times \frac{1}{50}$ from left to right. **Availability**

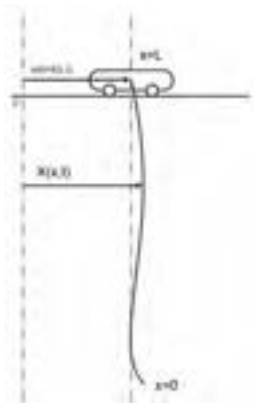
MEMS (microfab. resonators)

What can you achieve with these? (by combining them?)

Problem 1: propagation through a Kinematic Chain

Idea: deploy IMUs massively

Equip a mechanical system with **distributed** inertial sensors

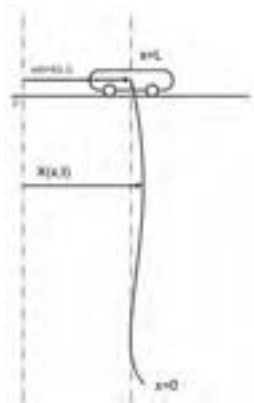


$$\text{heavy chain } \frac{\partial}{\partial x} \left(g x \frac{\partial X}{\partial x} \right) - \frac{\partial^2 X}{\partial t^2} = 0$$

Problem 1: propagation through a Kinematic Chain

Idea: deploy IMUs massively

Equip a mechanical system with distributed inertial sensors



$$\text{heavy chain } \frac{\partial}{\partial x} \left(g x \frac{\partial X}{\partial x} \right) - \frac{\partial^2 X}{\partial t^2} = 0$$

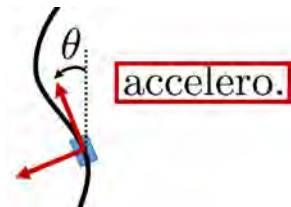
Accelerometer measurements

$$M = \begin{pmatrix} X(x, t) \\ x \end{pmatrix}, \quad \tan \theta = -\frac{\partial X}{\partial x}, \quad \frac{\partial}{\partial x} \left(g x \frac{\partial X}{\partial x} \right) - \frac{\partial^2 X}{\partial t^2} = 0$$

$$\begin{aligned} \text{accelero.} &= R_{\theta+\pi/2}^T \left(\ddot{M} + \begin{pmatrix} 0 \\ g \end{pmatrix} \right) \\ &\approx \begin{pmatrix} g \\ -\frac{\partial^2 X}{\partial t^2} + g \frac{\partial X}{\partial x} \end{pmatrix} \quad (\text{small angles}) \\ &= \begin{pmatrix} g \\ -g x \frac{\partial^2 X}{\partial x^2} \end{pmatrix} = \dots \begin{pmatrix} g \\ -\frac{g}{\tau'(x)} \tau(x) \frac{\partial^2 X}{\partial x^2} \end{pmatrix} \end{aligned}$$

$$\text{gyro.} = \dot{\theta} = -\frac{\partial^2 X}{\partial x \partial t}$$

$$\text{strain g.} = -\frac{\partial^2 X}{\partial x^2}$$



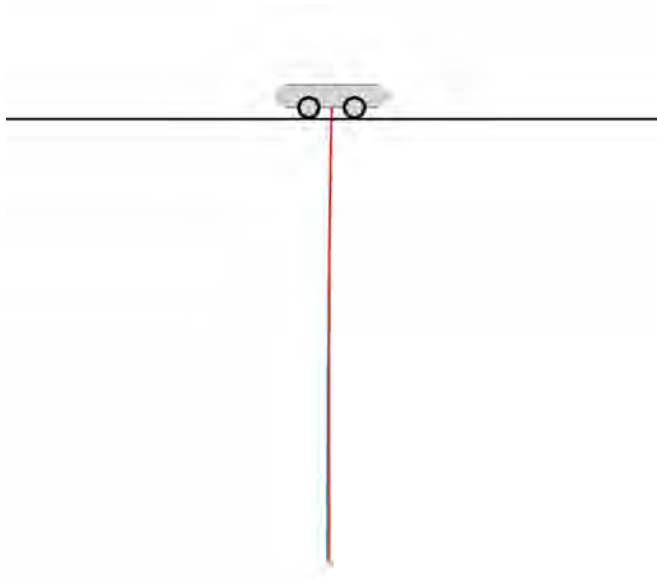
State reconstruction from accelerometer measurements

$$\text{accelero.}(x, t) = -gx \frac{\partial^2 X}{\partial x^2}(x, t)$$

Summation of distributed accelero.

$$\begin{aligned}\frac{\partial X}{\partial x}(x, t) &= \frac{\partial X}{\partial x}(L, t) + \int_x^L \frac{\text{accelero.}(\ell, t)}{g\ell} d\ell \\ X(x, t) &= X(L, t) - \int_x^L \frac{\partial X}{\partial x}(\ell, t) d\ell\end{aligned}$$

Simulation



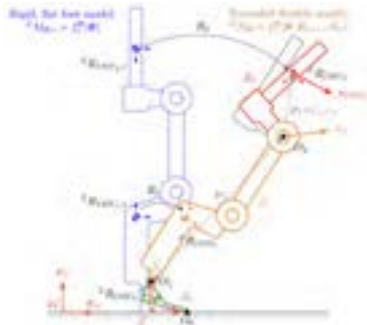
The real case study: Wandercraft exoskeleton

- Gyroscope
- Magnetometer
- Accelerometer



NS, LP, FdM, NP, +

Flexibilities and walk



Flexibilities

Reference frames

Positions or rigid bodies (local reference frame)

Kinematic chain from stance foot to swing foot

Not a standard robot: carries a non instrumented **patient**

System flexes, **foot deviates**, ground contact occurs too early, wrong mass transfer, ...

For any point M in the kinematic chain

$$M = f_{flex.}(\text{encoders}, R_0, R_1, \dots, R_n)$$

Measurement equation (not only gravity)

$$\begin{cases} \mathbf{y}_{a,i} = R_{IMU,i}^T (\ddot{\mathbf{p}}_{IMU,i} - \mathbf{g}) \\ \mathbf{y}_{g,i} = \mathbf{w}_{IMU,i}, \quad \dot{R}_{IMU,i} = R_{IMU,i} [\mathbf{w}_{IMU}]_{\times} \end{cases}$$

Velocity of each IMU (in its own reference frame)

$$\begin{aligned} \mathbf{v}_{IMU,i} &= R_{IMU,i}^T \dot{\mathbf{p}}_{IMU,i} \\ \dot{\mathbf{v}}_{IMU,i} &= -[\mathbf{y}_{g,i}]_{\times} \mathbf{v}_{IMU,i} + \mathbf{y}_{a,i} + R_{IMU,i}^T \mathbf{g} \end{aligned}$$

Approximation through a chain rule (from the stand foot)

$$\mathbf{v}^{\text{approx.}} : \begin{cases} \mathbf{v}_{O_i} = \mathcal{P}_{i-1} \mathbf{v}_{O_{i-1}} + \text{compo.}(\mathbf{y}_g, \text{encoders}) \times \ell_i \\ \mathbf{v}_{IMU,i} = \mathcal{Q}_i \mathbf{v}_{O_i} + [\mathbf{y}_g]_{\times} \ell_{IMU_i/O_i} \end{cases}$$

Observer (cont.)

Tilt vector (only relevant information, reduction of unknowns)

$$\mathbf{t} = R^T \mathbf{e}_z$$

For each rigid body in the kinematic chain

$$\begin{cases} \dot{\mathbf{v}} = -[\mathbf{y}_g]_{\times} \mathbf{v} + \mathbf{y}_a + g \mathbf{t} \\ \dot{\mathbf{t}} = -[\mathbf{y}_g]_{\times} \mathbf{t} \end{cases}$$

Cascade calculus from the ground to the swing foot

$$\begin{cases} \dot{\hat{\mathbf{v}}} = -[\mathbf{y}_g]_{\times} \hat{\mathbf{v}} + \mathbf{y}_a + g \hat{\mathbf{t}} + \alpha(\hat{\mathbf{v}} - \mathbf{v}^{\text{approx.}}) \\ \dot{\hat{\mathbf{t}}} = -[\mathbf{y}_g - \beta[\hat{\mathbf{t}}]_{\times} (\hat{\mathbf{v}} - \mathbf{v}^{\text{approx.}})]_{\times} \hat{\mathbf{t}} \end{cases}$$

$\alpha > \sqrt{\beta g}$, practical convergence (exponential, residual error, high gain)

$$\hat{R} = \exp \left[\arccos(\mathbf{e}_z^T \hat{\mathbf{t}}) \frac{\mathbf{e}_z \times \hat{\mathbf{t}}}{\|\mathbf{e}_z \times \hat{\mathbf{t}}\|} \right]_{\times} R_{\text{rigid}}$$

Experiments

Estimate and compensate flexibility-induced foot deviations



Robotics everywhere?



The missing ingredient: force control (not only position)
“Stick-slip” is prevalent, 3D, varying

Not only alter system behavior but **identify its parameters**

An example of advanced metrology: Ballistic wind tunnels

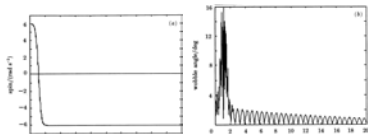
Before, let us discuss familiar complex mechanical systems

Celtic stone (rattleback)



(handcrafted by Pierre)

Reverses direction of rotation



¹ Garcia, Hubbard "Spin Reversal of the Rattleback: Theory and Experiment". Proc. of the Royal Society A (1988)

Routh sphere

Dynamically non-symmetric, unbalanced (displaced center of mass) ball rolling without slipping on a plane

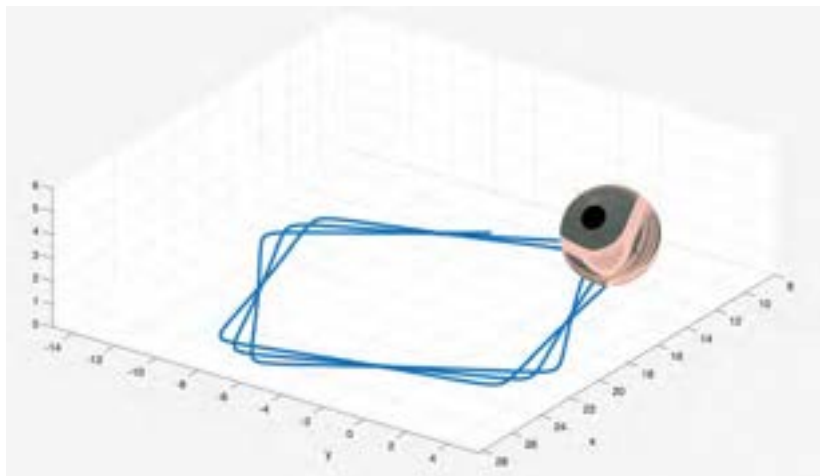
¹ E. J. Routh, The Advanced Part of A Treatise on the Dynamics of a System of Rigid Bodies, 6th ed., Macmillan, London, 1884



Trajectory

Non uniform motion is expected

Simulated motion



Model and derivation of sensors measurement

Euler equations

$$\mathbf{J} = \mathbf{I}_G + m \left(\|\mathbf{r}\|^2 \mathbf{E} - \mathbf{r} \mathbf{r}^T \right)$$

$$\boldsymbol{\omega} = \mathbf{J}^{-1} \mathbf{M}$$

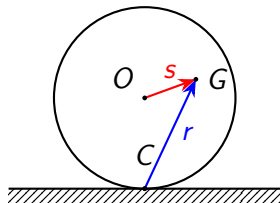
$$\dot{\mathbf{M}} = \mathbf{M} \times \boldsymbol{\omega} + m \left(\boldsymbol{\omega}^T (\mathbf{s} \times \mathbf{r}) \right) \boldsymbol{\omega} + m \frac{g}{a} \mathbf{r} \times \mathbf{s}$$

$$\dot{\mathbf{r}} = (\mathbf{r} - \mathbf{s}) \times \boldsymbol{\omega}$$

$$\dot{\mathbf{R}} = \mathbf{R} [\boldsymbol{\omega}]_{\times}$$

$$\dot{\mathbf{p}} = -a \mathbf{e}_z \times (\mathbf{R} \boldsymbol{\omega})$$

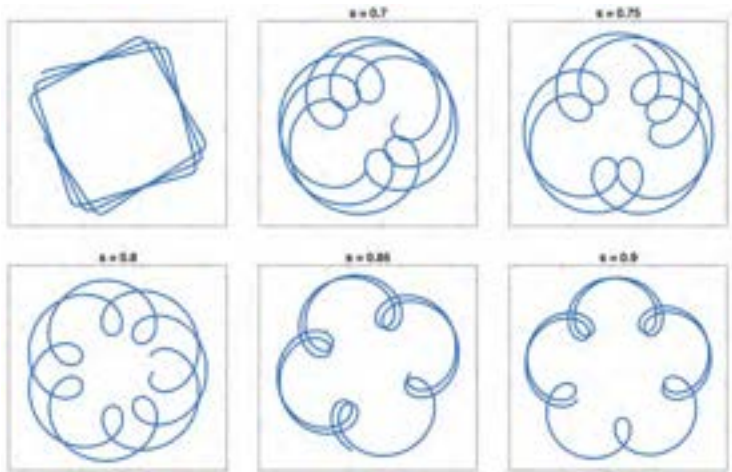
$\mathbf{s} \neq \mathbf{0}$: orientation dependent-torque,
non-holonomic constraint (increases
apparent inertia),



Spirograph



Spirograph when s is varied



Oscillations in the sensors

Gibbs-Appell equation of motion for the non-holonomic constraints

Routhian

$$R = \frac{1}{2}(\ddot{X}^2 + \ddot{Y}^2 + \ddot{Z}^2) + \frac{1}{2} \sum_i l_i \alpha_i^2$$

Intermediate frame $\mathcal{I}$, $(\mathbf{e}_1^i, \mathbf{e}_2^i, \mathbf{e}_3^i)$ (precession, tilt but no spin), angular velocity is $(\omega_1^i, \omega_2^i, \omega_3^i)^T$

$$\omega_1^i \mathbf{e}_1^i + \omega_2^i \mathbf{e}_2^i + (\omega_3^i - \dot{\psi}) \mathbf{e}_3^i = \dot{\theta} \mathbf{e}_1^i + \dot{\psi} \sin \theta \mathbf{e}_2^i + \dot{\phi} \cos \theta \mathbf{e}_3^i$$

Absolute **angular acceleration** vector is $\alpha = (\alpha_1, \alpha_2, \alpha_3)$ with

$$\alpha_1 = \dot{\omega}_1^i - (\omega_2^i)^2 \cot \theta + \dot{\omega}_2^i \dot{\omega}_3^i$$

$$\alpha_2 = \dot{\omega}_2^i + \omega_1^i \omega_2^i \cot \theta - \omega_1^i \dot{\omega}_3^i$$

$$\alpha_3 = \dot{\omega}_3^i$$

What does the nutation frequency tell us?

Non-holonomic constraints (rolling without slipping) give

$$\dot{X} = (s\omega_2^i - \omega_3^i) \cos \psi + \omega_1^i \sin \psi (\|\mathbf{s}\| \cos \theta + \sin \theta), \quad \dot{Y} = \dots$$

Tilt vector in the vicinity of the Euler angle **singularity** $\theta = 0$

Gibbs-Appell equation gives for the nutation frequency

$$A(s) \lambda^2 - K \lambda + H(s) = 0, \quad \omega_{\text{nut}} = \frac{\lambda_+(s) - \lambda_-(s)}{2}$$

with $A = I_1 + ms^2$, $K = (I_3 + ma(a + s))\omega_3(0)$, $H = mgs$

Want to know the value of s of your Routh sphere?

Just **spin it fast enough**, then watch it **roll** and measure the **frequency** of any sensor attached to it.

The real case study: Ballistic wind tunnel (high velocity)

- Gyroscope
- Magnetometer
- Accelerometer

ExoMars probe



Another version



Obtain attitude from magnetometer only

- 1 If **one angle and one direction** are known, the complete attitude can be estimated
- 2 If the **velocity** is known, the pitch angle can be deduced (via aerodynamics and a gain-scheduled observer).

① Attitude estimation

Projectile problem

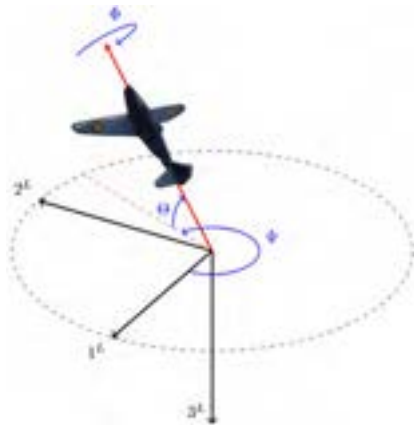
Only **one** direction and **one** angle are measured

Tait-Bryan seq. (ZYX e.g.)

yaw: Ψ

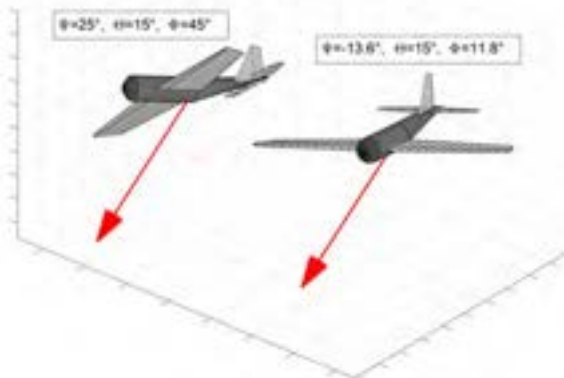
pitch: Θ

roll: Φ



Ambiguity of measurements (b_0) and pitch angles (15°) are identical

$$b_0 = (0.7412 \quad 0.8741 \quad 0.6671)^T$$
$$Y = (0.8865 \quad 0.4163 \quad 0.7983)^T$$



Quaternion analysis

Set of unit quaternions

$$\mathbb{Q} = \{q = (s, v \triangleq \text{dir}(q)) \in \mathbb{R} \times \mathbb{R}^3, |q| = 1\}, \text{ with } \otimes$$

$$q_1 \otimes q_2 = \begin{pmatrix} s_1 s_2 - v_1^T v_2 \\ s_1 v_2 + s_2 v_1 + v_1 \times v_2 \end{pmatrix}$$

neutral element $(1, 0, 0, 0)$, $q^{-1} = (s, -v)$, $q \otimes q^{-1} = (1, 0, 0, 0)$

Consider

$$R = F(q) := I_3 + 2s[v_\times] + 2[v_\times]^2 \quad (1)$$

which is “two-to-one” $\mathbb{Q} \rightarrow SO_3$

Measurements and Tait-Bryan angles

Direction measurement

$$Y = \text{dir} (q^{-1} \otimes \mathbf{p}(b_0) \otimes q) \triangleq q^{-1} b_0 q$$

with $\mathbf{p}(b_0) \triangleq (0, b_0)$ “pure” quaternion.

$$q \triangleq (q_1, q_2, q_3, q_4) \mapsto \begin{cases} T(q) = q_1 q_3 - q_2 q_4 \\ T_2(q) = q_1 q_4 + q_2 q_3, T_3(q) = 1 - 2(q_2^2 + q_3^2) \\ T_4(q) = q_1 q_2 + q_3 q_4, T_5(q) = 1 - 2(q_2^2 + q_3^2) \end{cases}$$

$$\Phi(q) = \arctan \frac{2T_4(q)}{T_5(q)}, \quad \Psi(q) = \arctan \frac{2T_2(q)}{T_3(q)}$$

Pitch measurement

$$\Theta(q) = \arcsin 2T(q)$$

Problem summary

Statement

Estimate (in real time) q from $Y = q^{-1}b_0q$ and $T(q)$.

$$\dot{R} = R[w_x], \quad \dot{q} = \frac{1}{2}q \otimes p(\omega) \quad (2)$$

Solution

$$\begin{cases} \dot{\hat{q}} = \frac{1}{2}\hat{q} \otimes p(\omega + k_p \sigma), & \sigma \triangleq Y \times (\hat{q}^{-1}b_0 \hat{q}) \\ \dot{\chi} = k_c \frac{T(q) - T(\bar{q})}{f(\nabla T(\bar{q})^T(p(b_0) \otimes \bar{q}))} \\ \bar{q} = \left(\cos \frac{\chi}{2} + \sin \frac{\chi}{2} b_0\right) \otimes \hat{q} \end{cases} \quad (3)$$

$$f : x \mapsto \sqrt{1+x^2} / \tanh(x)$$

$\bar{q}(t)$ converges towards $q(t)$ (conv. to be defined) when $t \rightarrow +\infty$

Sketch of proof

$$\dot{\hat{q}} = \frac{1}{2} \hat{q} \otimes \mathbf{p} (\omega + k_p \sigma)$$

vector error

$$\sigma \triangleq Y \times (\hat{q}^{-1} b_0 \hat{q}) = (q^{-1} b_0 q) \times (\hat{q}^{-1} b_0 \hat{q})$$

Some calculus yields (exponentially, the direction error is zero)

$$\lim_{t \rightarrow +\infty} \sigma(t) = 0$$

Extension by integrator on the gradient of pitch error to drive the system through the set $\mathcal{C}(q) \triangleq \{\hat{q}, \sigma(q, \hat{q}) = 0\}$, using an additional (output) rotation

$$\begin{cases} \dot{\chi} = k_c \frac{T(q) - T(\bar{q})}{f (\nabla T(\bar{q})^T (\mathbf{p}(b_0) \otimes \bar{q}))} \\ \bar{q} = \left(\cos \frac{\chi}{2} + \sin \frac{\chi}{2} b_0 \right) \otimes \hat{q} \end{cases}$$

Convergence set

$C(q)$ is the intersection of $\text{span}\{q, \mathbf{p}(b_0) \otimes q\}$ and the unit sphere of $\mathbb{R}^4$, i.e. a **circle of $\mathbb{R}^4$** . It can be parameterized as

$$C(q) = \{q_\ell \triangleq (\cos \ell) q + (\sin \ell) (\mathbf{p}(b_0) \otimes q), \ell \in [0, 2\pi]\} \quad (4)$$

Intersection of $C(q)$ with $T(q_\ell) = T(q)$ gives (with $b_0 \triangleq [a \ b \ c]^T$)

$$\begin{aligned} T(q_\ell) = T(q) = & (\cos^2 \ell) T(q) + (\cos \ell \sin \ell) (b T_3(q) - 2a T_2(q)) \\ & + (\sin^2 \ell) (-ac T_3(q) - 2bc T_2(q)) \\ & + (\sin^2 \ell) (c^2 - a^2 - b^2) T(q) \end{aligned}$$

which has **2 isolated solutions** $\ell \in [0, \pi[$: $\ell = 0$ and

$$\ell_{\#} = \frac{\pi}{2} + \arctan \left(\frac{b T_3(q) - 2a T_2(q)}{2(a^2 + b^2) T(q) + ac T_3(q) + 2bc T_2(q)} \right)$$

$q \neq q_{\#}$

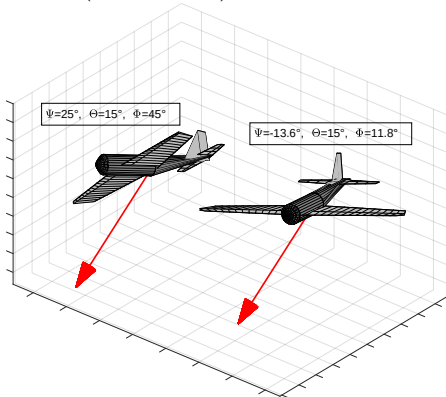
(and $\ell = \pi$ and $\ell = \ell_{\#} + \pi$, corresponding to $-q$ and $-q_{\#}$ (yielding identical rotations in SO_3).

Proposition

There exist *4 elements in $C(q)$* possessing the same pitch angle as q . These 4 are q , $-q$, $q_{\#}$ and $-q_{\#}$ yielding *two rotations*.

$$\mathbf{b}_0 = (0.7412 \quad 0.0741 \quad 0.6671)^T$$

$$\mathbf{Y} = (0.5065 \quad 0.4103 \quad 0.7583)^T$$



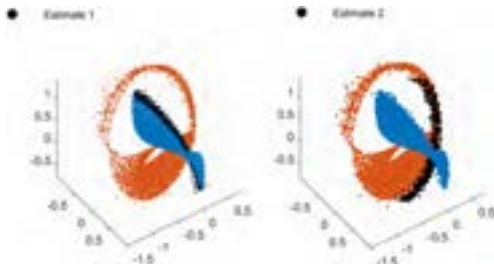
Theorem (cv. towards one of the two possible attitudes)

Let $\hat{q}(0)$ and $\chi(0)$ s.t. $\bar{q}(0) \notin E_0$. There exists $\eta_2 > 0$ s.t., for any $\varepsilon > 0$ small enough, there exist $\eta_1 > 0$ and $k_c^* > 0$ s.t., if $\|\sigma(0)\| < \eta_1$, $|T(q(0)) - T(\bar{q}(0))| < \eta_2$ and $k_c > k_c^*$, then there exists $K, \lambda > 0$ s.t. for any $t \geq 0$, one has

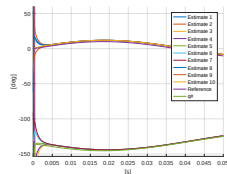
$$\min(|\delta(q, \bar{q})|, |\delta(-q, \bar{q})|, |\delta(q_\#, \bar{q})|, |\delta(-q_\#, \bar{q})|) < \varepsilon + K \exp^{-\lambda t}.$$

$\delta(., .)$: angle of relative quaternion

Practical behavior



Stereo. view of q , $q^\#$
and estimate $\bar{q}$ for various initial conditions.



Yaw angle

Straight-forward clarification

$\sigma(t)$ goes to zero at $t \rightarrow \infty$. Then $\bar{q}$ remains on side of the separating **hyperplane**. The projectile does not fly backwards.

Particular case

Pitch angle and Earth's magnetic field enable differentiation.
(except if the flight happens in a pure "North-South" plane).

② Estimating the velocity: equation of “exterior ballistics”

Euler-Poinsot equations (free rotation) with torques (Magnus $\mathcal{M}$, overturning $\mathcal{O}$, drag $\mathcal{D}$)

$$\begin{cases} I_l \dot{p} = -\mathcal{D}_p v p \\ I_t \dot{q} = (I_l - I_t) p r + \mathcal{M} v \beta + \mathcal{O} v^2 \alpha - \mathcal{D} v q \\ I_t \dot{r} = (I_l - I_t) p q + \mathcal{M} v \alpha + \mathcal{O} v^2 \beta - \mathcal{D} v r \end{cases} \quad (5)$$

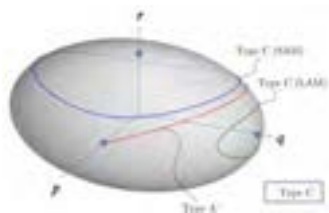
Type C reduction (intermediate frame),
parametric oscillator

$$\ddot{\xi} + v(H - iP)\dot{\xi} - v^2(M + iPT)\xi = -iPG$$

H, M, P, T, G : aero.

$$\xi(t) = A_n e^{i\omega_n t} + A_p e^{i\omega_p t} + A_0$$

Roll, Precession (rot. about roll axis),
Nutation (tilt of roll axis)



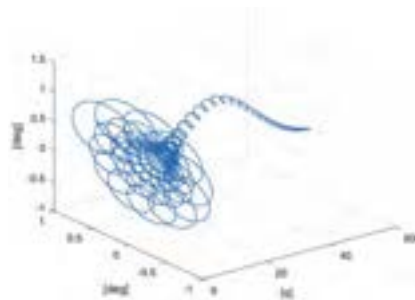
Angle of attack (aerodynamics)

Body attached **epicyclic** motion

$$\alpha(t) + i\beta(t) = -i(A_n e^{i(p-w_n)t} + A_p e^{i(p-w_p)t} + A_0 e^{ipt})$$

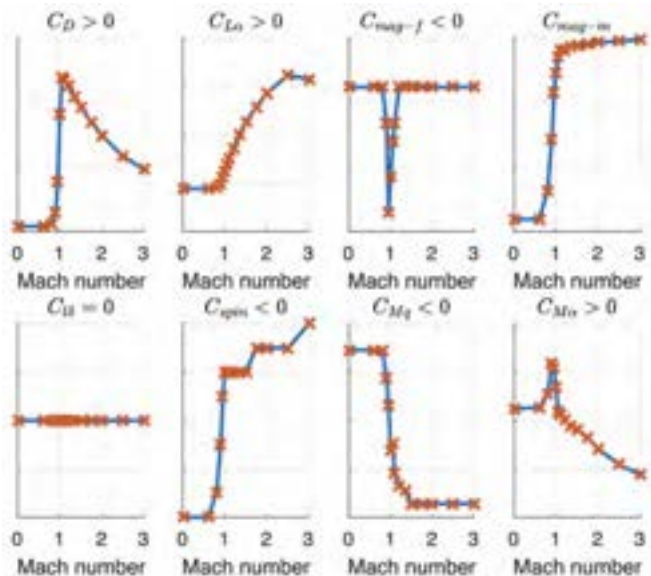
$$w_n = p \frac{I_l}{2I_t} + \frac{\sqrt{P_1^2 + P_2^2}}{2} \cos\left(\frac{1}{2} \arctan\left(\frac{P_2}{P_1}\right)\right)$$

$$w_p = p \frac{I_l}{2I_t} - \frac{\sqrt{P_1^2 + P_2^2}}{2} \cos\left(\frac{1}{2} \arctan\left(\frac{P_2}{P_1}\right)\right)$$

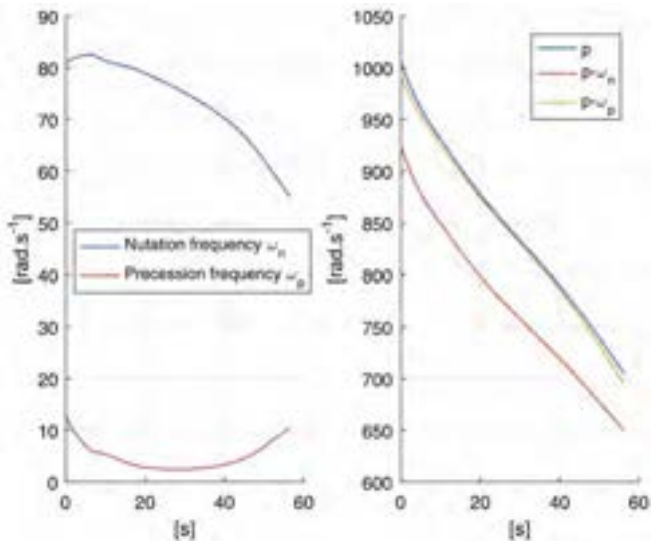


Analytic formula from aerodynamic forces and torques reference tables (drag, Magnus, lift).

Aerodynamic reference tables

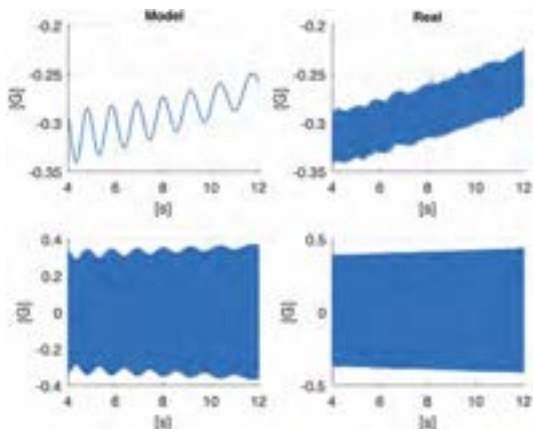


Theoretical frequencies



Magnetometer measurements

- Eddy currents
- Centrifugal effects
- Crosstalks
- Mis-alignments



“Vector space” frequency detection

$$s(t) = \sum_{k=1}^p \alpha_k \exp(i(w_k t + \phi_k)) + n(t)$$

$$y(t) = [s(t), \dots, s(t+m-1)]^T$$

$$C = E(y(t)y^*(t)) \in \mathbb{C}^{m \times m} = APA^* + \sigma^2 I_m$$

$$P = \text{diag}(\alpha_1^2, \dots, \alpha_p^2), \quad A = (a(w_1), \dots, a(w_p)) \in \mathbb{C}^{m \times p}$$

$$a(w) = (1, \exp(iw), \dots, \exp(iw(m-1)))^T$$

$$\text{eig}(C) : \lambda_1 \geq \dots \geq \lambda_p > \lambda_{p+1} = \dots = \lambda_m = \sigma^2 : G = [g_1, \dots, g_{m-p}]$$

Discriminating algebraic function

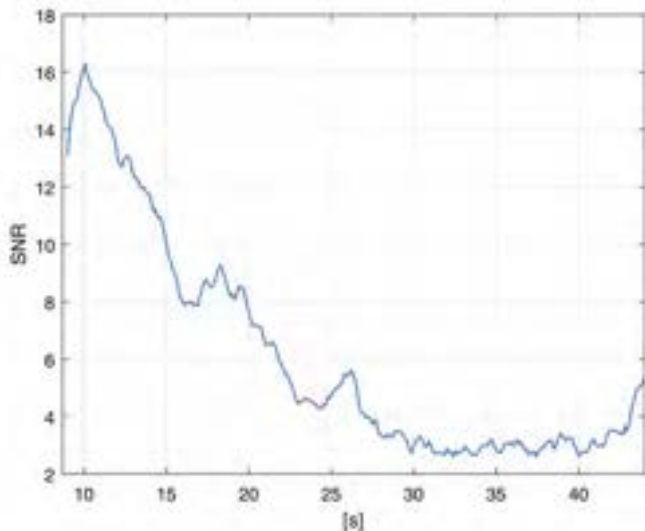
$$h_{\text{MU}}(w) = \frac{1}{\sum_{i=1}^{m-p} \|g_i^* a(w)\|^2} = \frac{1}{a^2(w) G G^* a(w)}$$

Detection of p peaks

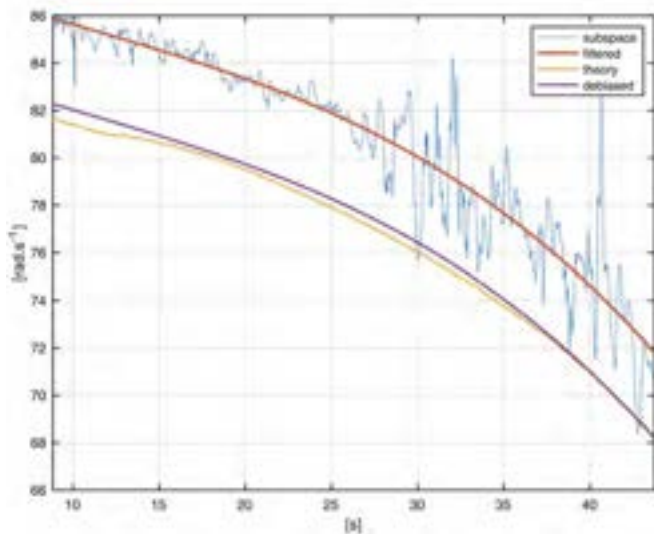
$$a^*(w_i) G G^* a(w_i) = 0, \quad i = 1, \dots, p$$

Recursive, embeddable implementation

Application: direct fire 40 s, 300 m/s, 900 rad/s

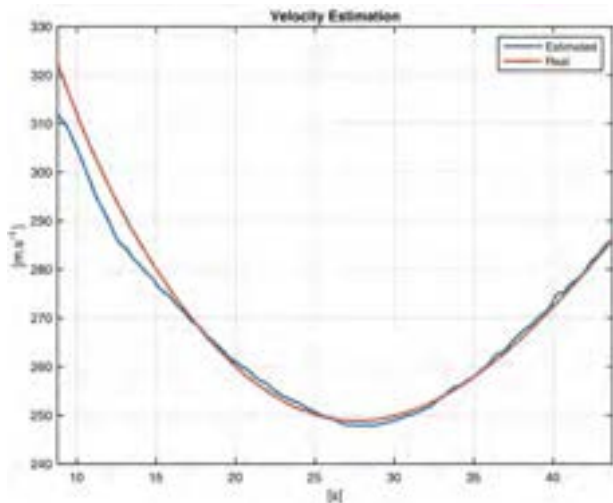


Dected frequencies



after roll-deconvolution

Velocity estimation



First velocity estimation without radar, GPS, or vision

Thank you for your attention